

Zum idealen Bose-Gas:

$$\sum_{n_\lambda=0}^{\infty} e^{-\beta(\epsilon_\lambda - \mu)n_\lambda} = \frac{1}{1 - e^{-\beta(\epsilon_\lambda - \mu)}}$$

nur für $\boxed{\epsilon_\lambda > \mu} \Rightarrow \beta(\epsilon_\lambda - \mu) > 0$

$$Z_G = \sum_{\{n_\lambda\}} e^{-\beta \sum_\lambda n_\lambda (\epsilon_\lambda - \mu)} \Rightarrow W_{\{n_\lambda\}} = \frac{1}{Z_G} e^{-\beta \sum_\lambda n_\lambda (\epsilon_\lambda - \mu)}$$

$$\begin{aligned} W_{\lambda_1}(n_{\lambda_1}) &= \sum_{\{n_\lambda, \lambda \neq \lambda_1\}} \frac{1}{\prod_\lambda Z_\lambda} \prod_{\lambda \neq \lambda_1} Z_\lambda e^{-\beta n_{\lambda_1} (\epsilon_{\lambda_1} - \mu)} \\ &= \frac{1}{Z_{\lambda_1}} e^{-\beta n_{\lambda_1} (\epsilon_{\lambda_1} - \mu)} \end{aligned}$$

$$\begin{aligned} \langle n_\lambda \rangle &= \frac{1}{Z_\lambda} \sum_{n_\lambda} n_\lambda e^{-n_\lambda x} & x &:= \beta(\epsilon_\lambda - \mu) \\ &= (1 - e^{-x})^{-1} (-1) \frac{\partial}{\partial x} \sum e^{-n_\lambda x} & Z_\lambda &= \frac{1}{1 - e^{-x}} \\ &= (e^{-x} - 1) \frac{\partial}{\partial x} \left(\frac{1}{1 - e^{-x}} \right) = (e^{-x} - 1) (-1) \frac{1}{(1 - e^{-x})^2} e^{-x} \\ &= \frac{e^{-x}}{1 - e^{-x}} = \boxed{\frac{1}{e^x - 1}} \quad \text{Bose-Funktion} \end{aligned}$$

(Bild auf Folie)

$$n_B = \frac{1}{e^x - 1}$$

$$x \ll 1 \quad n_B = \frac{1}{1 + x + \dots - 1} \approx \frac{1}{x}$$

$$x \gg 1 \quad n_B \approx e^{-x} \quad \stackrel{!}{=} \text{Maxwell Boltzmann-Gas}$$

Für hohe Temp ist $\mu \ll 0$

$$\text{so dass } \frac{1}{k_B T} (\epsilon_\lambda - \mu) \gg 1$$

Fermi-Gas : Pauli Verbot bedeutet $n_\lambda \in \{0; 1\}$

$$z_\lambda = 1 + e^{-x}$$

$$w_\lambda(n_\lambda) = \frac{1}{z_\lambda} e^{-n_\lambda x}$$

$$w_\lambda(n_\lambda=0) = \frac{1}{1 + e^{-x}}$$

$$w_\lambda(n_\lambda=1) = \frac{e^{-x}}{1 + e^{-x}}$$

$$\langle n_\lambda \rangle = 1 \cdot \frac{e^{-x}}{1 + e^{-x}} = \frac{1}{e^x + 1}$$

← "war" - " für Bose-Teilchen

Fermi-Funktion

$$x \ll -1 \Rightarrow n_F(\epsilon) \approx 1$$

$$\Leftrightarrow \epsilon_\lambda - \mu \ll -k_B T$$

$$x \gg 1 \Rightarrow n_F(\epsilon) \approx 0$$

für $T \rightarrow 0$ sind Zustände entweder besetzt oder

$$\text{leer. } n_F = \Theta(\mu - \epsilon)$$

(Bild auf Folie)

Ideales Fermi-Gas

$$\langle N \rangle = - \left. \frac{d\Omega}{d\mu} \right|_{T, V} = - \frac{d}{d\mu} (-k_B T) \sum_\lambda \ln(1 + e^{-\beta(\epsilon_\lambda - \mu)})$$

$$= k_B T \sum_\lambda \frac{1}{1 + e^{-\beta(\epsilon_\lambda - \mu)}} \beta e^{-\beta(\epsilon_\lambda - \mu)} = \sum_\lambda \langle n_\lambda \rangle$$

Kanonische Logik vs. großkanonische Logik

$$N \rightarrow \mu$$

(Folie)

$$\mu \rightarrow \langle N \rangle$$

$$S = - \left. \frac{d\Omega}{dT} \right|_{V, \mu} = \frac{d}{dT} \left(k_B T \sum_{\lambda} \ln(1 + e^{-\beta(\epsilon_{\lambda} - \mu)}) \right)$$

$$= k_B \sum_{\lambda} \ln() + k_B T \sum_{\lambda} \frac{1}{1 + e^{-\beta(\epsilon_{\lambda} - \mu)}} (\lambda)(\epsilon_{\lambda} - \mu) \left(\underbrace{\frac{1}{k_B T}}_{\beta} \right) e^{-x}$$

$$x_{\lambda} = \beta(\epsilon_{\lambda} - \mu)$$

$$= k_B \sum_{\lambda} \ln(1 + e^{-x_{\lambda}}) + \frac{x_{\lambda} e^{-x_{\lambda}}}{1 + e^{-x_{\lambda}}}$$

$$n_{\mp}(\epsilon_{\lambda}) = \frac{1}{e^{x_{\lambda}} + 1} \quad \left| \quad 1 - n_{\mp} = 1 - \frac{1}{e^{x_{\lambda}} + 1} = \frac{e^{x_{\lambda}}}{e^{x_{\lambda}} + 1} = \frac{1}{1 + e^{-x_{\lambda}}}\right.$$

$$= k_B \sum_{\lambda} -\ln[1 - n_{\mp}] + (1 - \langle n_{\lambda} \rangle) \ln \frac{1 - \langle n_{\lambda} \rangle}{\langle n_{\lambda} \rangle}$$

$$x_{\lambda} = \ln e^{x_{\lambda}}$$

$$\text{mit } e^{x_{\lambda}} + 1 = \frac{1}{n_{\mp}}$$

$$e^{x_{\lambda}} = \frac{1}{n_{\mp}} - 1 = \frac{1 - n_{\mp}}{n_{\mp}}$$

U auf Folie

