

# Hohlraum-Strahlung (Photonen, Phononen)

$$\sum_{\vec{s}} f(\epsilon) = 2V \int \frac{d^3p}{(2\pi\hbar)^3} f(\epsilon) = 2V \int \frac{4\pi p^2 dp}{(2\pi\hbar)^3} f(\epsilon)$$

$$\epsilon = cp \quad dp = \frac{d\epsilon}{c} \quad \Rightarrow \gamma(\epsilon) = \frac{4\pi\epsilon^2}{(2\pi\hbar c)^3} \quad \text{für Photonen}$$

$$f_{\text{ges}} (2s+1) = 2 \quad ?$$

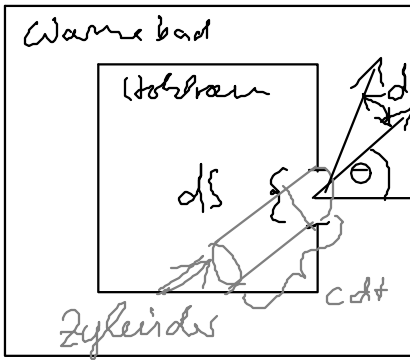
$Z(T, V)$  kanonisch oder  $Z(T, V, \mu=0)$  großkanonisch

$$x = \beta \hbar \omega$$

$$\Omega = 2V k_B T \int_0^\infty d\omega \frac{4\pi\omega^2}{(2\pi c)^3} \ln(1 - e^{-x}) = \frac{2V k_B T}{(\beta \hbar)^3} \int_0^\infty dx \frac{4\pi x^2}{(2\pi c)^3} \ln(1 - e^{-x})$$

$$= \frac{2V (k_B T)^4}{(2\pi \hbar c)^3} 4\pi \int_0^\infty dx x^2 \ln(1 - e^{-x})$$

$$U = F + TS = F + T \left( 4 \frac{F}{T} \right) = -3F$$



$$C_V = T \left( \frac{\partial S}{\partial T} \right)_V \quad \text{Wärmekapazität}$$

$$U d\omega = \frac{1}{V} \frac{1}{e^{\beta \hbar \omega} - 1} \frac{4\pi k^2 d\epsilon}{(2\pi)^3} (2V) \hbar \omega$$

$$\hbar \omega \ll k_B T : u(\omega, T) = \frac{\hbar}{\pi^2 c^3} \frac{\omega^3}{1 + \beta \hbar \omega - 1}$$

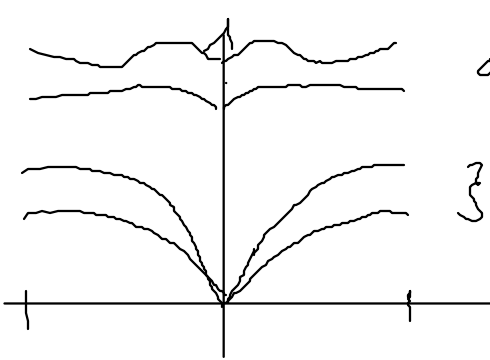
$$dV_{\text{Zylinder}} = c dt (d\ell \cos \theta)$$

$$dE = dV u(\omega, T) d\omega d\Omega$$

Abgestrahlte Leistung

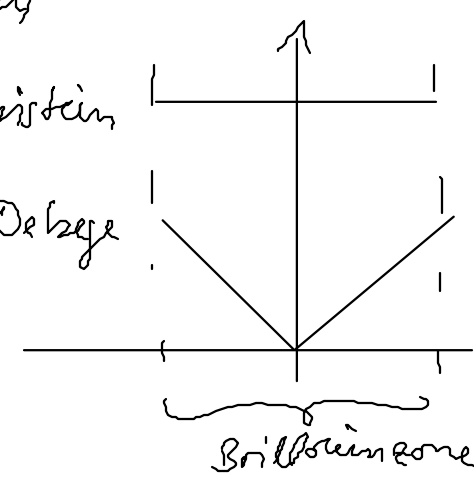
$$dI(\theta) = \frac{dN}{dt} E$$

**für Phononen** (nur akustische Phononen)



optische  
} akustische

Einstein  
Debye



# ideale Systeme

- Wenn exakte Berechnung möglich  $\Rightarrow$  Näherungen verwenden

z.B.  $V_{\text{Coulomb}} = \frac{e^2}{4\pi\epsilon_0 |\mathbf{r}_1 - \mathbf{r}_2|}$  (cgs-System)

$$H = H_{\text{kin}}(\{p_i\}) + H_{\text{pot}}(\{r_i\}) + H_{\text{int}}(\{r_i\})$$

$\underbrace{\hspace{10em}}_{\text{Wechselwirkung}}$

- Gelas Beobachtung
- gut reale Systeme

## Virialentwicklung

$$Z = \text{Tr}(e^{-\beta H})$$

- Virialkoeffizienten B und C

$$f(x) = \ln(1-x) \quad f'(x) = -\frac{1}{1-x} \quad f'' = -\frac{1}{(1-x)^2} \quad f''' = -\frac{2}{(1-x)^3}$$

$$f^{(n)} = -\frac{(n-1)!}{(1-x)^n}$$

$$f(x) = \sum_{n=0}^{\infty} \frac{1}{n!} x^n f^{(n)}(x=0) = -\sum_{n=1}^{\infty} \frac{x^n}{n}$$

$$-\Omega = k_B T Z_1 e^{\beta\mu} + k_B T \left( Z_2 - \frac{Z_1^2}{2} \right) e^{2\beta\mu}$$

$$N = Z_1 e^{\beta\mu} + 2 \left( Z_2 - \frac{Z_1^2}{2} \right) e^{2\beta\mu}$$

$$e^{2\beta\mu} = \left( \frac{N}{Z_1} \right)^2$$

$$-\Omega = k_B T \left( Z_1 e^{\beta\mu} + \left( Z_2 - \frac{Z_1^2}{2} \right) e^{2\beta\mu} \right) = k_B T \left( N - \left( Z_2 - \frac{Z_1^2}{2} \right) e^{2\beta\mu} \right)$$

$$\bar{R} = \frac{1}{N} \frac{1}{2} \frac{1}{\lambda^3}$$

$$r = r_1 - r_2$$

$$\int d^3r_1 d^3r_2 f(|r_1 - r_2|) = \int d^3R \int d^3r f(|r|)$$

$$B = -V \left( \frac{\frac{1}{2} V \int d^3r e^{-\beta V(r)}}{\lambda^3 \left( \frac{V}{\lambda^3} \right)^2} - \frac{1}{2} \right) = -\frac{1}{2} V \left( \frac{\int e^{-\beta V(r)} d^3r}{V} - 1 \right)$$

$$= -\frac{1}{2} \int d^3r (e^{-\beta V(r)} - 1)$$

$$b = \frac{1}{2} \frac{4\pi}{3} (2r_0)^3 \quad B = b - a/\rho$$

↑ 2 Teilchen (ununterscheidbar)

$$P = k_B T N \left( 1 + b n - \frac{a n^2}{k_B T} \right)$$

$$P V + N a n = k_B T N (1 + b n)$$

$$V(P + a n^2) = k_B T N (1 + b n)$$

$$V(P + a n^2)(1 - b n) = k_B T N$$

$$\frac{N}{V} = n$$

$$\frac{1}{1+x} \approx 1-x$$

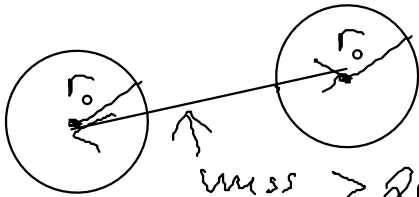
für

$x \ll 1$

$n b \ll 1$

$n a \ll 1$

Van der Waals-Gleichung



mus  $> 2r_0$  sein, sonst ausgeschlossenes Volumen  
(Teilchen ineinander, nicht  
mehr frei)

$$V - V_0 = V - Nb$$

$$(V - V_0)^n = (V - Nb)^n = V^n \left( 1 - \frac{Nb}{V} \right)^n$$

$$N! = \left( \frac{N}{e} \right)^N$$